

Families of periodic solutions for some Hamiltonian PDEs (with G. Arioli)

- (1) The problem etc.
- (2) Main results
- (3) Numerical results
- (4) Proofs

We consider **time-periodic solutions** for the **nonlinear wave equation** ($\mu = 1$) and the **nonlinear beam equation** ($\mu = 2$)

$$\partial_t^2 u(t, x) + (-1)^\mu \partial_x^{2\mu} u(t, x) = f(u(t, x)), \quad (t, x) \in \mathbb{R} \times (0, \pi),$$

with Dirichlet BCs. These PDEs are **Hamiltonian** with

$$H(u, v) = \int_0^\pi \left[\frac{1}{2} (\partial_x^\mu u)^2 + \frac{1}{2} v^2 - F(u) \right] dx, \quad F' = f.$$

From a period 2π one can get “related” periods via scaling. Changes f unless homogeneous.

Our motivation:

- Observed instabilities in a bridge model [Arioli, Gazzola 2000].
- CAP for Hamiltonian and/or parabolic PDEs with potential small denominator issues.

Existing relayed work:

Variational methods for period 2π and related:

$\mu = 1$ [Rabinowitz 1978; Rabinowitz 1981; ...]

$\mu = 2$ [Lee 2000; Liu 2002; Liu 2004; ...]

Perturbative methods for small u and positive-measure sets of periods near “special” values:

$\mu = 1$ [Berti 2007; Gentile, Mastropietro, Procesi 2005; Gentile, Procesi 2009]

$\mu = 2$ [Mastropietro, Procesi 2006; Gentile, Procesi 2009]

We restrict to $f(\mathbf{u}) = \sigma \mathbf{u}^3$ with $\sigma = \pm 1$.

Setting $\mathbf{u}(t, x) = u(\alpha t, x)$, where $\frac{2\pi}{\alpha}$ is the desired period for \mathbf{u} , we arrive at the equation

$$L_\alpha u = \sigma u^3, \quad L_\alpha = \alpha^2 \partial_t^2 + (-1)^\mu \partial_x^{2\mu},$$

where $u = u(t, x)$ is 2π -periodic in t and satisfies Dirichlet boundary conditions at $x = 0, \pi$.

$\mu = 1$: Nonlinear **wave equation** $\alpha^2 \partial_t^2 u - \partial_x^2 u = \sigma u^3$.

$\mu = 2$: Nonlinear **beam equation** $\alpha^2 \partial_t^2 u + \partial_x^4 u = \sigma u^3$.

Consider the vector space \mathcal{A}° of all real analytic functions

$$u = \sum_{n,k} u_{n,k} P_{n,k}, \quad P_{n,k}(t, x) = \cos(nt) \sin(kx).$$

We restrict our analysis to the subspace \mathcal{B} consisting of all $u \in \mathcal{A}^\circ$ with the property that $u_{n,k} \neq 0$ only if n and k are both odd.

Notice that

$$L_\alpha P_{n,k} = \lambda_{n,k} P_{n,k}, \quad \lambda_{n,k} = k^{2\mu} - (\alpha n)^2 = (k^\mu + \alpha n)(k^\mu - \alpha n).$$

We only consider α values for which $\lambda_{n,k} \neq 0$ for all odd n and k .

This includes the set \mathbb{Q}_o of rationals $\alpha = p/q$ with p and q of opposite parity.

Definition. A solution $u \in \mathcal{B}$ of the equation $L_\alpha u = \sigma u^3$ will be called a *type (1, 1) solution* if $|u_{n,k}| < |u_{1,1}|$ whenever $n > 1$ or $k > 1$.

First consider the **nonlinear wave equation** for some rational values of α . Our sample set:

$$Q_1 = \left\{ \frac{3}{8}, \frac{5}{12}, \frac{7}{16}, \frac{9}{20}, \frac{13}{28}, \frac{1}{2}, \frac{15}{28}, \frac{11}{20}, \frac{9}{16}, \frac{7}{12}, \frac{5}{8}, \frac{9}{14}, \frac{11}{16}, \frac{7}{10}, \frac{13}{18}, \frac{3}{4}, \frac{11}{14}, \frac{5}{6}, \frac{7}{8}, \frac{9}{10}, \frac{11}{12}, \frac{13}{14}, \frac{17}{18} \right\}.$$

Theorem 1. For each $\alpha \in Q_1$ the equation $\alpha^2 \partial_t^2 u - \partial_x^2 u = u^3$ has a solution $u \in \mathcal{B}$ of type (1, 1) with $|u_{1,1}| > \sqrt{2(1-\alpha)}$.

Remark. Every solution $u \in \mathcal{B}$ of the equation $\alpha^2 \partial_t^2 u - \partial_x^2 u = u^3$ with $\alpha \in Q_0$ yields a solution $\tilde{u} \in \mathcal{B}$ of the equation $\alpha^2 \partial_t^2 \tilde{u} - \partial_x^2 \tilde{u} = -\tilde{u}^3$, and vice-versa. The functions u and \tilde{u} are related via $\tilde{u}(t, x) = \alpha^{-1} u(x - \pi/2, t - \pi/2)$.

Next we consider irrational values of α .

Unfortunately we have to switch to the **nonlinear beam equation**.

Still difficult to construct non-small solutions for specific α .

Best known: $\alpha = 1/\sqrt{c}$ where c is an integer that is not the square of an integer.

By Siegel's theorem on integral points on algebraic curves of genus one,

$$c\lambda_{n,k} = ck^4 - n^2 \rightarrow \infty \quad \text{as} \quad n \vee k \rightarrow \infty.$$

Unfortunately we have no useful bounds ...

So we make an assumption:

Theorem 2. *Let $\alpha = 1/\sqrt{3}$. Assume that $|3k^4 - n^2| \geq 39$ for all $k \geq 9$ and all $n \in \mathbb{N}$. Then the equation $\alpha^2 \partial_t^2 u + \partial_x^4 u = u^3$ has a solution $u \in \mathcal{B}$ of type $(1, 1)$ with $|u_{1,1}| > 1$.*

We have verified the assumption $\min_n |3k^4 - n^2| \geq 39$ for $9 \leq k \leq 10^{12}$.

Our third result concerns irrational values of α that are close to the rationals in

$$Q_2 = \left\{ \frac{1}{4}, \frac{3}{10}, \frac{9}{20}, \frac{1}{2}, \frac{7}{12}, \frac{5}{8}, \frac{3}{4}, \frac{5}{6}, \frac{7}{6}, \frac{5}{4}, \frac{19}{14}, \frac{17}{12}, \frac{31}{20}, \frac{13}{8}, \frac{31}{18}, \frac{61}{34} \right\}.$$

Theorem 3. *For each $r \in Q_2$ there exists a set $R \subset \mathbb{R}$ of positive measure that includes r as a Lebesgue density point, such that for each $\alpha \in R$, the equation $\alpha^2 \partial_t^2 u + \partial_x^4 u = \sigma u^3$ with $\sigma = \text{sign}(1 - \alpha)$ has a solution $u \in \mathcal{B}$ of type $(1, 1)$ with $|u_{1,1}| > \sqrt{2|1 - \alpha|}$.*

Remark. In all of the equations considered, other types of solutions can be obtained via scaling: If $u \in \mathcal{A}^\circ$ satisfies the equation $L_\alpha u = \sigma u^3$, and if we define

$$\tilde{u}(t, x) = b^\mu u(at, bx), \quad \tilde{\alpha} = \alpha b^\mu / a, \quad (\star)$$

with b and a nonzero integers, then \tilde{u} belongs to \mathcal{A}° and satisfies $L_{\tilde{\alpha}} \tilde{u} = \sigma \tilde{u}^3$.

In our proofs we solve $L_\alpha u = \sigma u^3$ via the fixed point equation

$$u = \mathcal{F}_\alpha(u) \stackrel{\text{def}}{=} L_\alpha^{-1} \sigma u^3, \quad \sigma = \text{sign}(1 - \alpha).$$

For **numerical experiments** we use Fourier polynomials

$$u = \sum_{\substack{n \leq N \\ k \leq K}} u_{n,k} P_{n,k}, \quad P_{n,k}(t, x) = \cos(nt) \sin(kx).$$

and **truncate** u^3 to wavenumbers $n \leq N$ and $k \leq K$.

As $N \rightarrow \infty$ the equation becomes Hamiltonian, even if $K < \infty$.

Definition for the $K < \infty$ equation. *The union of all smooth branches that include a solution of type $(1, 1)$ will be referred to as the **(1, 1) branch**. Scaling each solution on the $(1, 1)$ branch via (\star) yields what we will call the **(a, b) branch**.*

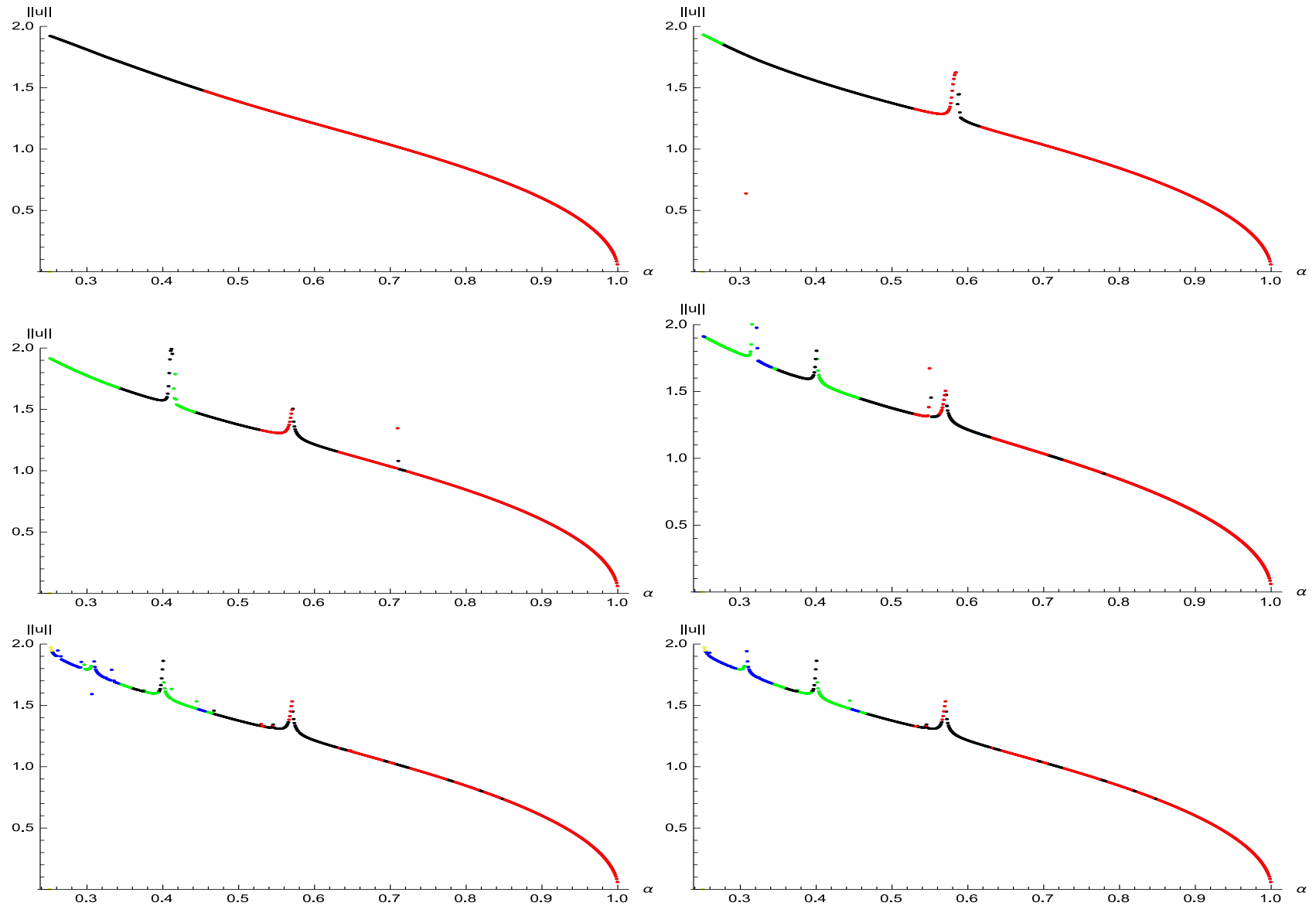
In the following **graphs** we show the norm

$$\|u\|_0 = \sum_{n,k} |u_{n,k}|$$

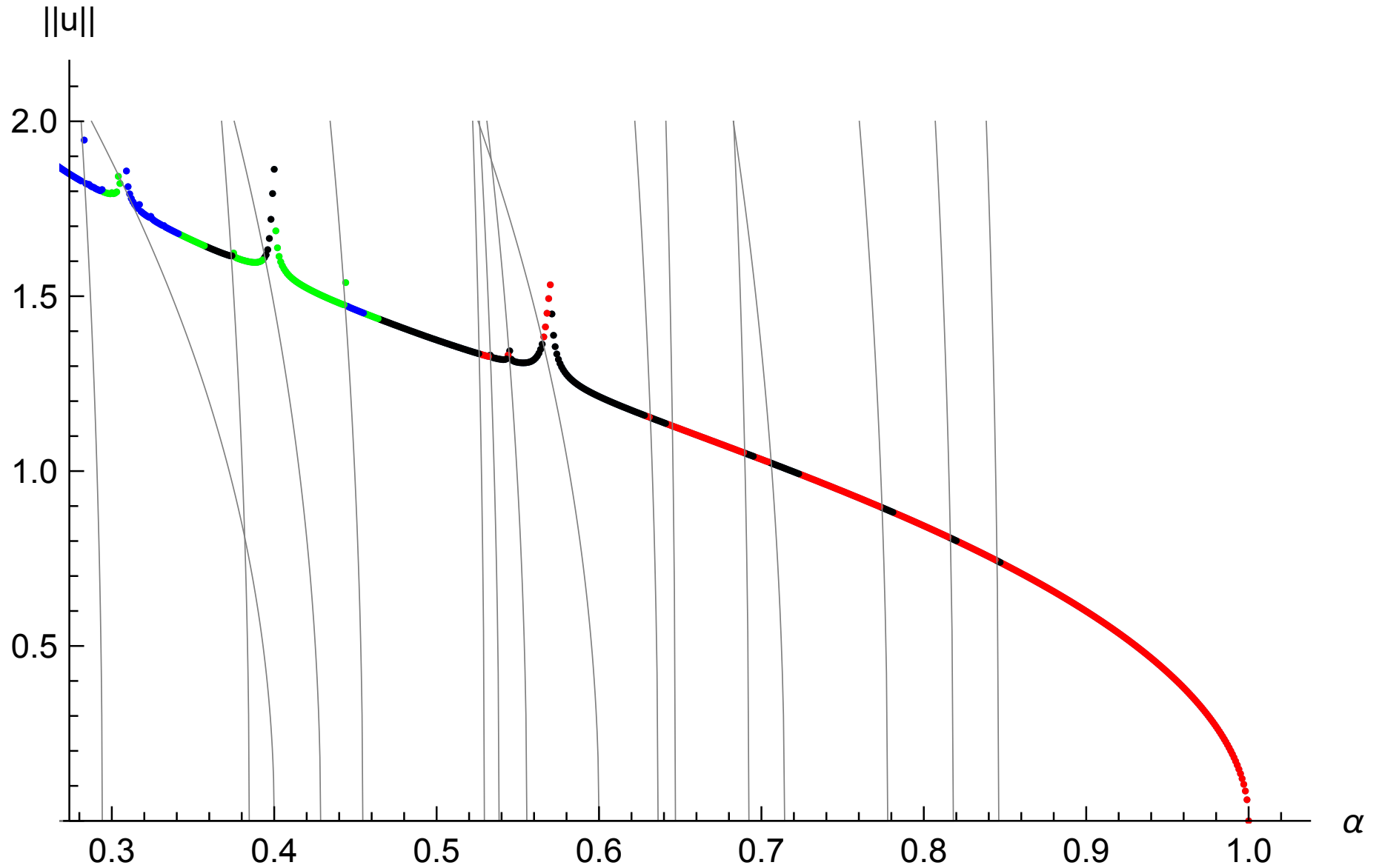
of the numerical solution u as a function of α .

C o l o r s encode the index of u : the number of eigenvalues larger than 1 of $D\mathcal{F}_\alpha(u)$.

The (1,1) branch for the nonlinear wave equation $\alpha^2 u_{tt} - u_{xx} = u^3$ truncated at $N = K = 3, 5, 7, 9, 19, 39$.

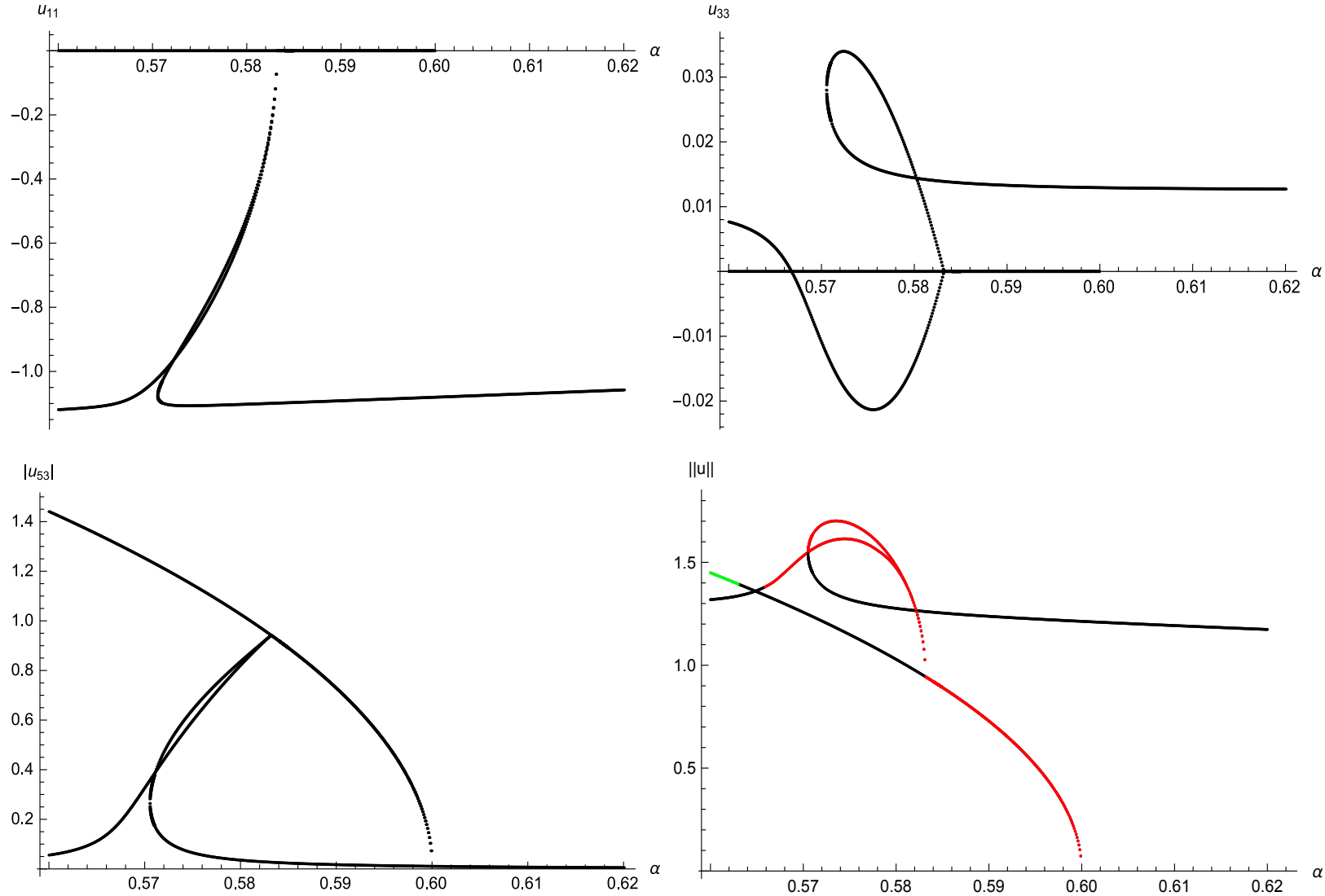


The (1,1) branch for the truncated nonlinear wave equation $\alpha^2 u_{tt} - u_{xx} = u^3$ and some other (a,b) branches (thin lines).

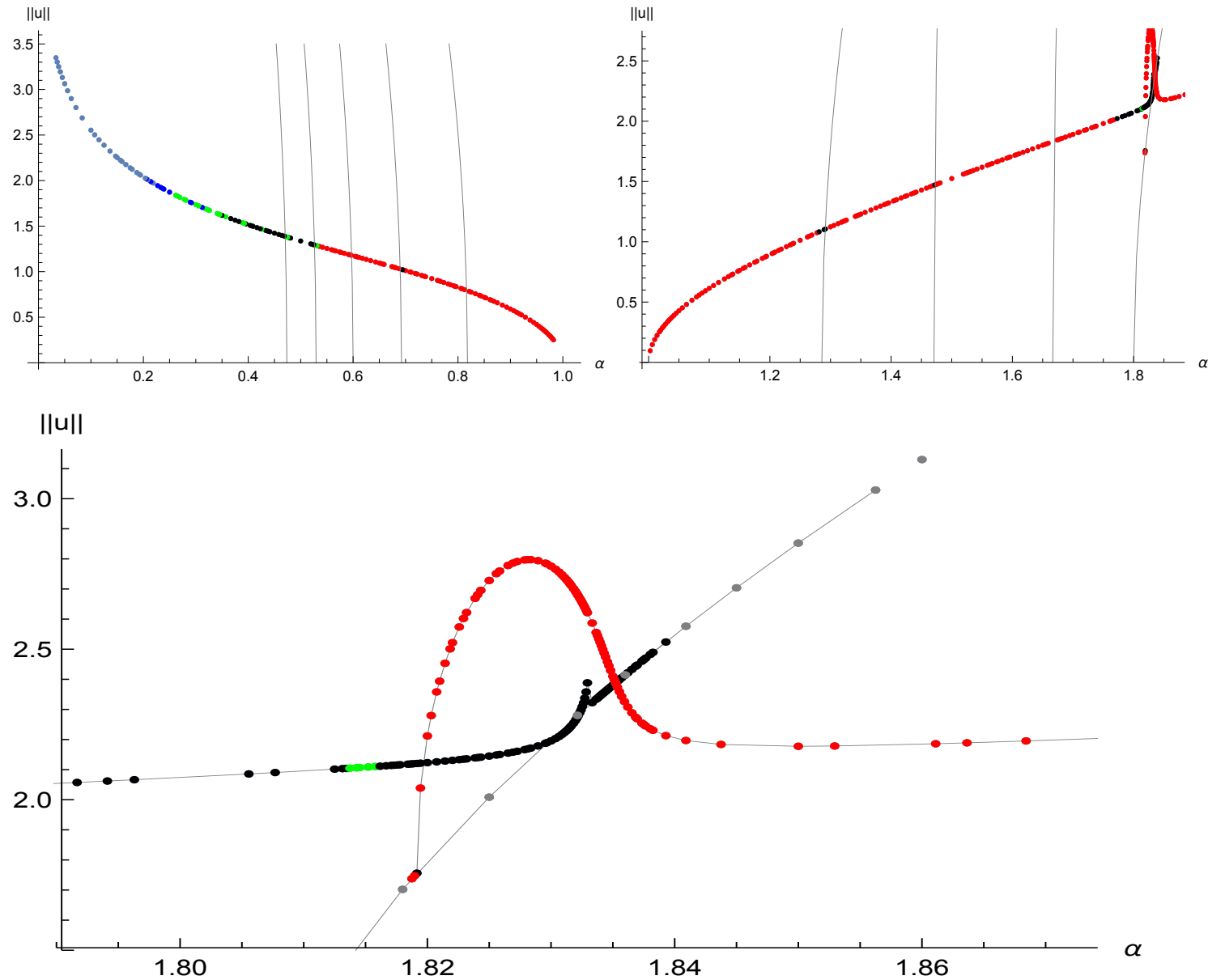


The **nonlinear wave equation**, truncated at $N \gg K = 7$.

The (1,1) branch undergoes a fold bifurcation at $\alpha \simeq 0.571$
and a pitchfork bifurcation involving the (5,3) branch at $\alpha \simeq 0.585$.



The **nonlinear beam equation** $\alpha^2 u_{tt} + u_{xxxx} = \pm u^3$ truncated at $K = 63$ and $N = 127$.



The proofs of Theorems 1,2,3 use the contraction mapping theorem.

Given $\rho = (\rho_1, \rho_2)$ with $\rho_j > 0$ denote by \mathcal{A}_ρ° the closure of \dots with respect to the norm

$$\|u\|_\rho = \sum_{n,k} |u_{n,k}| \varrho_1^n \varrho_2^k, \quad \varrho_j = 1 + \rho_j.$$

Let $\mathcal{B}_\rho = \mathcal{B} \cap \mathcal{A}_\rho^\circ$. Consider a quasi-Newton map associated with \mathcal{F}_α ,

$$\mathcal{N}_\alpha(h) = \mathcal{F}_\alpha(u_0 + Ah) - u_0 + (I - A)h,$$

where u_0 is an approximate fixed point and A an approximate inverse of $I - D\mathcal{F}_\alpha(u_0)$.

Denote by B_δ the open ball of radius δ in \mathcal{B}_ρ , centered at the origin.

Theorem 3 is proved by verifying the following bounds.

Lemma 4 *For each $r \in Q_2$ there exists a set $R \subset \mathbb{R}$ of positive measure that includes r as a Lebesgue density point, a pair ρ of positive real numbers, a Fourier polynomial $u_0 \in \mathcal{B}_\rho$, a linear isomorphism $A : \mathcal{B}_\rho \rightarrow \mathcal{B}_\rho$, and positive constants K, δ, ε satisfying $\varepsilon + K\delta < \delta$, such that for every $\alpha \in R$ the map \mathcal{N}_α defined as above is analytic on B_δ and satisfies*

$$\|\mathcal{N}_\alpha(0)\|_\rho < \varepsilon, \quad \|D\mathcal{N}_\alpha(h)\|_\rho < K, \quad h \in B_\delta.$$

Compactness of $L_\alpha^{-1} : \mathcal{B}_\rho \rightarrow \mathcal{B}_\rho$.

Define $\lceil s \rceil = \text{dist}(s, \mathbb{Z})$. A simple estimate on the eigenvalues of L_α is

$$\beta^2 |\lambda_{n,k}| = (\beta k^\mu + n) |\beta k^\mu - n| \geq (2(\beta k^\mu \vee n) - \lceil \beta k^\mu \rceil) \lceil \beta k^\mu \rceil, \quad \beta = \alpha^{-1}.$$

If $\alpha = p/q$ with p odd and q even: $\lceil \beta k^\mu \rceil \geq 1/p$ for k odd; so in this case L_α^{-1} is compact.

For $\mu = 2$ and irrational α we can use the following.

Let $(\psi_1, \psi_2, \psi_3, \dots)$ be a summable sequence of nonnegative real numbers.

Proposition 5. *Let $m \geq 1$. Consider an interval J_m of length $m^{-2} \leq |J_m| \leq 1$. Then*

$$\{\beta \in J_m : \lceil \beta k^2 \rceil \geq \psi_k \quad \text{for all } k \geq m\}$$

has measure at least $(1 - 4\Psi_m)|J_m|$, where $\Psi_m = \sum_{k \geq m} \psi_k$.

Applying this with $|J_m| = 1$ and $\psi_k = k^{-3/2}$ yields the

Corollary 6. *For almost every $\alpha \in \mathbb{R}$ the operator $L_\alpha = \alpha^2 \partial_t^2 + \partial_x^4$ has a compact inverse.*

Subspaces for error terms: $u \in \mathcal{B}_{\rho,\nu,\kappa}$ iff $u \in \mathcal{B}_\rho$ and $u_{n,k} = 0$ whenever $n < \nu$ or $k < \kappa$.

Our enclosures for $u \in \mathcal{B}_\rho$ consist of interval enclosures for each $c_{n,k}$ and for the norm of each $E_{\nu,\kappa}$ in a representation

$$u = \sum_{\substack{n \leq N \\ k \leq K}} c_{n,k} P_{n,k} + \sum_{\substack{\nu \leq 2N \\ \kappa \leq 2K}} E_{\nu,\kappa}, \quad E_{\nu,\kappa} \in \mathcal{B}_{\rho,\nu,\kappa}.$$

Estimating the map $u \mapsto u^3$ on \mathcal{B}_ρ is “standard”.

The operator norm of $L_\alpha^{-1} : \mathcal{B}_{\rho,\nu,\kappa} \rightarrow \mathcal{B}_{\rho,\nu,\kappa}$ is bounded by $\beta^2 / \phi(\nu, \kappa)$ where

$$\phi(\nu, \kappa) = \inf_{\substack{n \geq \nu \\ k \geq \kappa}} \beta^2 |\lambda_{n,k}| = \inf_{\substack{n \geq \nu \\ k \geq \kappa}} (\beta k^\mu + n) |\beta k^\mu - n|, \quad \beta = \alpha^{-1}.$$

Here ν, κ, n, k are odd positive integers. To prove Lemma 4 we use

Lemma 7. *Let $r = p/q$ with p odd and q even. Given odd positive integers κ and ν , there exists a set $R \subset \mathbb{R}$ of positive measure that includes r as a Lebesgue density point, such that for all $\alpha \in R$,*

$$\phi(\nu, \kappa) \geq \frac{7}{4p} \left[(\beta \kappa^2 \vee \nu) - \frac{7}{16p} \right], \quad \beta = \alpha^{-1}.$$

Idea of the proof: In the above **inf** distinguish between $k \geq m$ and $k < m$.

For $k \geq m$ use Proposition 5 with J_m centered at r . And for $k < m$ use that α is close to r . Do this for increasingly large m .

Some references

- P.H. Rabinowitz, *Free vibration for a semilinear wave equation*, Comm. Pure Appl. Math.. **31**, 31–68 (1978).
- P.H. Rabinowitz, *On nontrivial solutions of a semilinear wave equation*, Ann. Scuola Norm. Sup. Pisa Cl. Sci. **8**, 647–657 (1981).
- M. Berti, *Nonlinear Oscillations of Hamiltonian PDEs*, Birkhäuser Verlag, 2007.
- W.M. Schmidt, *Metrical theorems on fractional parts of sequences*, Trans. Amer. Math. Soc. **110**, 493–518 (1964).
- L.H. Eliasson, B. Grébert, S. Kuksin, *KAM for the non-linear beam equation*, 2 Preprints.
- G. Gentile, V. Mastropietro, M. Procesi, *Periodic solutions for completely resonant nonlinear wave equations with Dirichlet boundary conditions*, Commun. Math. Phys. **256**, 437–490 (2005).
- V. Mastropietro, M. Procesi, *Lindstedt series for periodic solutions of beam equations with quadratic and velocity dependent nonlinearities*, Commun. Pure Appl. Anal. **5**, 1, 128 (2006).
- G. Gentile, M. Procesi, *Periodic solutions for a class of nonlinear partial differential equations in higher dimension*, Commun. Math. Phys. **289**, 863–906 (2009).
- C. Lee, *Periodic solutions of beam equations with symmetry*, Nonlin. Anal. T.M.A. **42**, 631–650 (2000)
- J.Q. Liu, *Free vibrations for an asymmetric beam equation*, Nonlin. Anal. T.M.A. **51**, 487–497 (2002)
- J.Q. Liu, *Free vibrations for an asymmetric beam equation, II*, Nonlin. Anal. T.M.A. **56**, 415–432 (2004)
- G. Arioli, F. Gazzola, *On a nonlinear nonlocal hyperbolic system modeling suspension bridges*, Milan J. Math. **83**, 211–236 (2015).
- G. Arioli, F. Gazzola, *Torsional instability in suspension bridges: the Tacoma Narrows Bridge case*, Preprint mp_arc 15-83.
- G. Arioli, H. Koch, *Non-symmetric low-index solutions for a symmetric boundary value problem*, J. Differ. Equations, **252** 448–458 (2012).
- G. Arioli, H. Koch, *Some symmetric boundary value problems and non-symmetric solutions*, J. Differ. Equations **259**, 796–816 (2015).
- G. Arioli, H. Koch, *Computer-Assisted Methods for the Study of Stationary Solutions in Dissipative Systems, Applied to the Kuramoto-Sivashinski Equation*, Arch. Rat. Mech. Anal. **197**, 1033–1051 (2010).
- G. Arioli, H. Koch, *Integration of Dissipative Partial Differential Equations: A Case Study*, SIAM J. Appl. Dyn. Syst. **9** 1119–1133 (2010).